NETWORK SCIENCE

## Destruction perfected

Pinpointing the nodes whose removal most effectively disrupts a network has become a lot easier with the development of an efficient algorithm. Potential applications might include cybersecurity and disease control. SEE LETTER P.65

ISTVÁN A. KOVÁCS & ALBERT-LÁSZLÓ BARABÁSI

n enduring truth of network science is that the removal of a few highly connected nodes, or hubs, can break up a complex network into many disconnected components<sup>1</sup>. Sometimes, a fragmented and inactive network is more desirable than a functioning one. Consider, for example, the need to eliminate bacteria by disrupting their molecular network or by vaccinating a few individuals in a population to break up the contact network through which a pathogen spreads. In a quest to find the silver bullets that can effectively dismantle large networks, Morone and Makse<sup>2</sup> (page 65 of this issue) have developed an algorithm that achieves this by identifying sets of network nodes known as influencers.

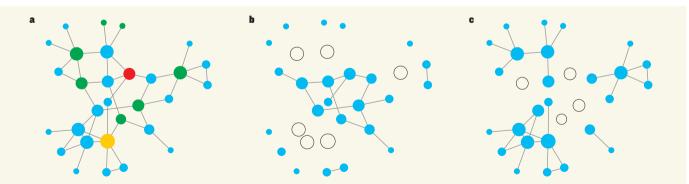
It is not certain whether targeting and removing network hubs — defined as the nodes with the largest number of links — can inflict maximum disruption on a network. It may be more effective to eliminate a combination of hubs and central, but less-well-connected, nodes. The removal of hubs is usually preferred because they are easy to locate, whereas identifying the optimal set of nodes for

which deletion would cause maximum damage is a non-deterministic polynomial-time hard (NP-hard) problem<sup>3</sup>. This means that it is computationally feasible only for small networks. Morone and Makse attack the problem of network disruption by mapping the integrity of a tree-like random network into optimal percolation<sup>4,5</sup> theory. From this, they derive an energy function with a minimum that corresponds to the set of nodes that need to be eliminated, to yield a network whose largest cluster is as small as possible. Although identifying this minimum is still an NPhard problem, the authors were inspired by the energy function's shape to find a simple algorithm that offers an approximate solution.

To do this, Morone and Makse introduce the concept of collective influence, which is the product of the node's reduced degree (the number of its links minus one) and the sum of the reduced degrees of the nodes that are a certain number of steps away from it (Fig. 1). Collective influence describes how many other nodes can be reached from a given node, assuming that nodes of high collective influence have a crucial role in the network. The collective-influence-based algorithm then sequentially removes nodes, starting with those that have the highest collective influence

(known as influencers) and recalculating the collective influence of the rest following each operation. The authors show that, for large networks, removing the set of influencers identified by this algorithm is more effective in fragmenting a network than removing the hubs, or than removing nodes that are identified through other algorithms, such as PageRank<sup>6</sup> or closeness centrality<sup>7</sup>. The set of influencers identified by the authors contains many nodes with few connections. This highlights the fact that the importance of a node in ensuring a network's integrity is determined not only by the number of direct links it has to other nodes, but also by which other nodes it is connected to.

The collective-influence algorithm is remarkable for its computational complexity because it requires only  $N^2 \log N$  computations to dismantle a network that contains N number of nodes. Its complexity is reduced to  $N \log N$  if, instead of individual nodes, a fixed fraction of the total is removed at each step of the computation. The authors compare their method to the predictions of spin-glass theory, which was originally developed to describe the properties of disordered magnets and has found a range of applications in network analysis. They conclude that the nodes prioritized



**Figure 1** | **Optimal network demolition.** Morone and Makse<sup>2</sup> introduce an algorithm that allows them to efficiently dismantle networks. The authors define the collective influence of a network node as the product of its reduced degree (the number of its nearest connections, k, minus one), and the total reduced degree of all nodes at distance d from it (defined as the number of steps from it). **a**, In this network, for d = 2, the red node with k = 4 has the highest collective influence, because the total reduced degree of the nodes at d = 2 from it (green and yellow circles) is 21. This yields a collective influence of  $3 \times 21 = 63$ . The most connected hub, with k = 6 (yellow circle), has a

collective influence of 60. **b**, Removing the 6 nodes with the highest k (white circles) causes considerable damage to the network, but leaves a sub-network that contains 12 nodes unperturbed. **c**, By contrast, the algorithm developed by the authors allows them to identify a set of nodes (known as influencers) according to their collective influence. Using this, the removal of four influencer nodes (white circles) results in a fragmented network in which the largest connected cluster that remains has only ten nodes. This illustrates the algorithm's effectiveness over conventional methods for prioritizing network destruction.

ISTVÁN A. KOVÁCS & ALBERT-LÁSZLÓ BARABÁS

by the collective-influence algorithm represent an approximate solution, which has a size close to that of the theoretical optimal solution. On the basis of spin-glass theory, we expect that the collective-influence solution has only a small overlap with the optimal solution, and hence must be treated with caution. However, the influencers found by collective influence are more effective in destroying a network than nodes selected by other methods. So even though the collective-influence method is approximate, it is faster and more efficient.

As with any new algorithm, open questions abound. The collective-influence algorithm has only one free parameter — the distance, expressed in the number of steps, from any given node. At zero distance, the collective influence of a node is equal to the square of its reduced degree, and so in this case the algorithm simply removes the hubs. To improve the algorithm's accuracy, one must choose a non-zero distance — but one that is not too large, because for large distances the boundaries of the network are reached, diminishing a node's collective influence (the collective influence approaches zero). Although Morone and Makse find that any distance greater than one works, a firm criterion for choosing an optimal value is lacking and would be desirable. Finally, because the authors designed their algorithm to work on networks that are locally tree-like, further work and quantitative evidence are needed on its expected accuracy for networks with loops, such as most social networks.

The collective-influence algorithm, just like similar algorithms, removes a node together with all its links. However, for many systems, node removal is too drastic an intervention. Softer touches, such as removing or rewiring specific links, are more tractable and desirable. For example, these approaches are relevant for networks in biological cells, in which many diseases are caused by mutations that result in deletion of links rather than the complete removal of nodes<sup>8</sup>. Understanding such 'edgetic' effects, and designing algorithms that can detect the minimum number of links to delete so as to achieve a given outcome, remains a challenge for future work.

The identification of optimal influencers, at either the node or the link level, is the first step towards building networks that would be robust against both attacks and failures. Mastering the design principles of such super-robust networks could have profound implications for anything from cybersecurity to the design of an attack- and error-tolerant power grid, and may even allow us to develop drugs that can rescue a cellular network from its diseased state with minimal side effects.

István A. Kovács and Albert-László Barabási are at the Center for Complex Network Research and in the Department of Physics, Northeastern University, Boston, Massachusetts 02115, USA. e-mail: alb@neu.edu

- Albert, R., Jeong, H. & Barabási, A.-L. Nature 406, 378–382 (2000).
- 2. Morone, F. & Makse, H. A. Nature 524, 65-68 (2015).
- 3. Garey, M. R. & Johnson, D. S. in Computers and Intractability: A Guide to the Theory of
- NP-completeness (Freeman, 1979).
- Hashimoto, K. Adv. Stud. Pure Math. 15, 211–280 (1989).
- Karrer, B., Newman, M. E. J. & Zdeborová, L. *Phys. Rev. Lett.* 113, 208702 (2014).
- Brin, S. & Page, L. Proc. 7th Int. World Wide Web Conf. 30, 107–117 (1998).
- 7. Freeman, L. C. Soc. Networks 1, 215–239 (1978–79).
- 3. Sahni, N. et al. Cell **161**, 647–660 (2015).

#### DIABETES

## A smart insulin patch

A microneedle-containing patch that is designed to sense elevated blood glucose levels and to respond by releasing insulin could offer people with diabetes a less-painful and more-reliable way to manage their condition.

#### OMID VEISEH & ROBERT LANGER

iabetes is widely recognized as one of the biggest medical challenges of the twenty-first century, afflicting more than 280 million people globally¹. People with diabetes must tirelessly self-monitor their blood glucose levels and inject the correct dose of the glucose-lowering hormone insulin to keep their blood glucose levels in the

Skin

Glucose

Glucose

Hypoxiaresposive
polymer

Figure 1 | A microneedle patch to monitor glucose and release insulin. Yu et al.4 have developed a smart insulin-releasing patch made of 121 nanoparticle-containing microneedles. The patch painlessly penetrates the interstitial fluid between subcutaneous skin cells. The nanoparticles in each needle contain insulin and the glucose-sensing enzyme glucose oxidase, which converts glucose to gluconic acid. These molecules are surrounded by a hypoxia-responsive polymer. Increases in glucose oxidase activity in response to glucose elevation produce a low-oxygen environment in the nanoparticles, which is sensed by the hypoxia-responsive polymer, triggering disassembly of the nanoparticles and the release of insulin.

normal range<sup>2</sup>. This treatment regime involves challenges — it requires painful and inconvenient subcutaneous injections, is imprecise, and can cause serious problems if insulin dosage is not closely tuned to the patient's immediate physiological needs<sup>3</sup>. Reporting in *Proceedings of the National Academy of Sciences*, Yu *et al.*<sup>4</sup> describe a glucose-responsive microneedle patch that can be painlessly applied to the skin and that releases insulin as blood glucose levels increase.

'Smart' glucose-responsive insulin-based therapies involve the automatic release of insulin in response to increases in blood glucose concentration. Smart therapies can improve disease control and limit the potential for excessively low blood glucose levels, which is a potentially deadly effect of excessive insulin dosing<sup>3</sup>. To mimic the physiological needs of a patient accurately, such therapies must respond rapidly to elevated glucose levels, and must release insulin with kinetics that closely mirror those of a healthy pancreas.

One type of smart therapy makes use of microcomputer-controlled insulin-delivery systems. These systems couple implantable continuous glucose monitors (CGMs) to automated pumps, and administer insulin through a subcutaneously inserted cannula tube. They are currently being evaluated in the clinic, and have shown promise in helping patients to achieve their target blood glucose level more regularly<sup>5,6</sup>. However, the sensors of current CGMs must be calibrated many times a day using hand-held glucometers. They produce blood-glucose measurements that lag behind true blood glucose levels by 5-15 minutes, hampering efforts to maintain a healthy range<sup>3</sup>. They are also the size of pagers, and the implanted sensors and cannula increase the risk of infection and require frequent maintenance and replacement to combat the body's immune response, increasing inconvenience, discomfort and cost to the patient<sup>3</sup>.

The microneedle-patch device developed by Yu and colleagues is a 6-millimetre-square



# Influence maximization in complex networks through optimal percolation

Flaviano Morone<sup>1</sup> & Hernán A. Makse<sup>1</sup>

The whole frame of interconnections in complex networks hinges on a specific set of structural nodes, much smaller than the total size, which, if activated, would cause the spread of information to the whole network<sup>1</sup>, or, if immunized, would prevent the diffusion of a large scale epidemic<sup>2,3</sup>. Localizing this optimal, that is, minimal, set of structural nodes, called influencers, is one of the most important problems in network science<sup>4,5</sup>. Despite the vast use of heuristic strategies to identify influential spreaders<sup>6-14</sup>, the problem remains unsolved. Here we map the problem onto optimal percolation in random networks to identify the minimal set of influencers, which arises by minimizing the energy of a many-body system, where the form of the interactions is fixed by the nonbacktracking matrix<sup>15</sup> of the network. Big data analyses reveal that the set of optimal influencers is much smaller than the one predicted by previous heuristic centralities. Remarkably, a large number of previously neglected weakly connected nodes emerges among the optimal influencers. These are topologically tagged as low-degree nodes surrounded by hierarchical coronas of hubs, and are uncovered only through the optimal collective interplay of all the influencers in the network. The present theoretical framework may hold a larger degree of universality, being applicable to other hard optimization problems exhibiting a continuous transition from a known phase<sup>16</sup>.

The optimal influence problem was initially introduced in the context of viral marketing<sup>1</sup>, and its solution was shown to be NP-hard<sup>4</sup> for a generic class of linear threshold models of information spreading<sup>17,18</sup>. Indeed, finding the optimal set of influencers is a many-body problem in which the topological interactions between them play a crucial role<sup>13,14</sup>. On the other hand, there has been an abundant production of heuristic rankings to identify influential nodes and 'superspreaders' in networks<sup>6–12,19</sup>. The main problem is that heuristic methods do not optimize a global function of influence. As a consequence, there is no guarantee of their performance.

Here we address the problem of quantifying nodes' influence by finding the optimal (that is, minimal) set of structural influencers. After defining a unified mathematical framework for both immunization and spreading, we provide its optimal solution in random networks by mapping the problem onto optimal percolation. In addition, we present CI (Collective Influence), a scalable algorithm to solve the optimization problem in large-scale real data sets. The thorough comparison with competing methods (Supplementary Information section I<sup>20</sup>) ultimately establishes the better performance of our algorithm. By taking into account collective influence effects, our optimization theory identifies a new class of strategic influencers, called 'weak nodes', which outrank the hubs in the network. Thus, the top influencers are highly counterintuitive: low-degree nodes play a major broker role in the network, and despite being weakly connected, can be powerful influencers.

The problem of finding the minimal set of activated nodes<sup>17,18</sup> to spread information to the whole network<sup>4</sup> or to optimally immunize a network against epidemics<sup>11</sup> can be exactly mapped onto optimal percolation (see Supplementary Information section IIB). This mapping

provides the mathematical support to the intuitive relation between influence and the concept of cohesion of a network: the most influential nodes are the ones forming the minimal set that guarantees a global connection of the network<sup>5,9,10</sup>. We call this minimal set the 'optimal influencers' of the network. At a general level, the optimal influence problem can be stated as follows: find the minimal set of nodes which, if removed, would break down the network into many disconnected pieces. The natural measure of influence is, therefore, the size of the largest (giant) connected component as the influencers are removed from the network.

We consider a network composed of N nodes tied with M links with an arbitrary-degree distribution. Let us suppose we remove a certain fraction q of the total number of nodes. It is well known from percolation theory<sup>21</sup> that, if we choose these nodes randomly, the network undergoes a structural collapse at a certain critical fraction where the probability of existence of the giant connected component vanishes, G = 0. The optimal influence problem corresponds to finding the minimum fraction  $q_c$  of influencers to fragment the network:  $q_c = \min\{q \in [0,1]: G(q) = 0\}$ .

Let the vector  $\mathbf{n} = (n_1, \dots, n_N)$  represent which node is removed  $(n_i = 0, \text{ influencer})$  or left  $(n_i = 1, \text{ the rest})$  in the network  $(q = 1 - 1/N \sum_i n_i)$ , and consider a link from i to j  $(i \rightarrow j)$ . The order parameter of the influence problem is the probability that i belongs to the giant component in a modified network where j is absent,  $v_{i \rightarrow j}$  (refs 22, 23). Clearly, in the absence of a giant component we find  $\{v_{i \rightarrow j} = 0\}$  for all  $i \rightarrow j$ . The stability of the solution  $\{v_{i \rightarrow j} = 0\}$  is controlled by the largest eigenvalue  $\lambda(\mathbf{n}; q)$  of the linear operator  $\hat{\mathcal{M}}$ ,

defined on the  $2M \times 2M$  directed edges as  $\mathcal{M}_{k \to \ell, i \to j} \equiv \frac{\partial v_{i \to j}}{\partial v_{k \to \ell}} \Big|_{\{v_{i \to j} = 0\}}$ . We find for locally tree-like random graphs (see Fig. 1a and Supplementary Information section II):

$$\mathcal{M}_{k \to \ell, i \to j} = n_i \mathcal{B}_{k \to \ell, i \to j} \tag{1}$$

where  $\mathcal{B}_{k \to \ell, i \to j}$  is the non-backtracking matrix of the network<sup>15,24</sup>. The matrix  $\mathcal{B}_{k \to \ell, i \to j}$  has non-zero entries only when  $(k \to \ell, i \to j)$  form a pair of consecutive non-backtracking directed edges, that is,  $(k \to \ell, \ell \to j)$  with  $k \neq j$ . In this case  $\mathcal{B}_{k \to \ell, \ell \to j} = 1$  (equation (13) in Supplementary Information). Powers of the matrix  $\hat{\mathcal{B}}$  count the number of non-backtracking walks of a given length in the network (Fig. 1b)<sup>24</sup>, much in the same way as powers of the adjacency matrix count the number of paths<sup>5</sup>. Operator  $\hat{\mathcal{B}}$  has recently received a lot of attention thanks to its high performance in the problem of community detection<sup>25,26</sup>. We show its topological power in the problem of optimal percolation.

Stability of the solution  $\{v_{i \to j} = 0\}$  requires  $\lambda(\mathbf{n};q) \le 1$ . The optimal influence problem for a given  $q \ge q_c$  can be rephrased as finding the optimal configuration  $\mathbf{n}$  that minimizes the largest eigenvalue  $\lambda(\mathbf{n};q)$  (Fig. 1c). The optimal set  $\mathbf{n}^*$  of  $Nq_c$  influencers is obtained when the minimum of the largest eigenvalue reaches the critical threshold:

$$\lambda(\mathbf{n}^*; q_c) = 1 \tag{2}$$

<sup>&</sup>lt;sup>1</sup>Levich Institute and Physics Department, City College of New York, New York, New York 10031, USA.

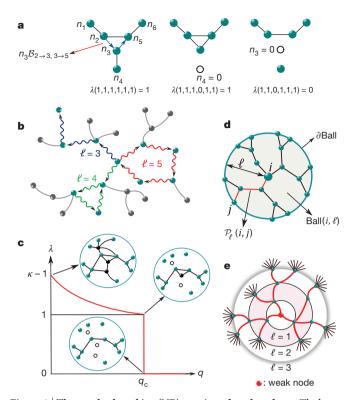


Figure 1 | The non-backtracking (NB) matrix and weak nodes. a, The largest eigenvalue  $\lambda$  of  $\mathcal{M}$  exemplified on a simple network. The optimal strategy for immunization and spreading minimizes  $\lambda$  by removing the minimum number of nodes (optimal influencers) that destroys all the loops. Left panel, the action of the matrix  $\mathcal{M}$  is on the directed edges of the network. The entry  $\mathcal{M}_{2\to3,3\to5} = n_3 \mathcal{B}_{2\to3,3\to5} = n_3$  encodes the occupancy  $(n_3 = 1)$  or vacancy  $(n_3 = 0)$  of node 3. In this particular case, the largest eigenvalue is  $\lambda = 1$ . Centre panel, non-optimal removal of a leaf,  $n_4 = 0$ , which does not decrease  $\lambda$ . Right panel, optimal removal of a loop,  $n_3 = 0$ , which decreases  $\lambda$  to zero. b, A NB walk is a random walk that is not allowed to return back along the edge that it just traversed. We show a NB open walk ( $\ell = 3$ ), a NB closed walk with a tail ( $\ell = 4$ ), and a NB closed walk with no tails ( $\ell = 5$ ). The NB walks are the building blocks of the diagrammatic expansion to calculate  $\lambda$ . c, Representation of the global minimum over n of the largest eigenvalue  $\lambda$  of  $\mathcal{M}$  versus q. When  $q \ge q_c$ , the minimum is at  $\lambda = 0$ . Then, G = 0 is stable (still, non-optimal configurations exist with  $\lambda > 1$  for which G > 0). When  $q < q_c$ , the minimum of the largest eigenvalue is always  $\lambda > 1$ , the solution  $\hat{G} = 0$  is unstable, and then G > 0. At the optimal percolation transition, the minimum is at  $\mathbf{n}^*$  with  $\lambda(\mathbf{n}^*, q_c) = 1$ . For q = 0, we find  $\lambda = \kappa - 1$  ( $\kappa = \langle k^2 \rangle / \langle k \rangle$ , where *k* is the node degree) which is the largest eigenvalue of  $\hat{\mathcal{B}}$  for random networks<sup>25</sup> with all nodes present ( $n_i = 1$ ). When  $\lambda = 1$ , the giant component is reduced to a tree plus one single loop (unicyclic graph), which is suddenly destroyed at the transition  $q_c$  to become a tree, causing the abrupt fall of  $\lambda$ to zero. **d**, Ball(i,  $\ell$ ) of radius  $\ell$  around node i is the set of nodes at distance  $\ell$ from i, and  $\partial$ Ball is the set of nodes on the boundary. The shortest path from i to j is shown in red. e, Example of a weak node: a node with a small number of connections surrounded by hierarchical coronas of hubs at different  $\ell$  levels.

The formal mathematical mapping of the optimal influence problem to the minimization of the largest eigenvalue of the modified nonbacktracking matrix for random networks, equation (2), represents our first main result.

An example of a non-optimized solution corresponds to choosing  $n_i$  at random and decoupled from the non-backtracking matrix<sup>23,27</sup> (random percolation<sup>21</sup>, Supplementary Information section IID). In the optimized case, we seek to derandomize the selection of the set  $n_i = 0$  and optimally choose them to find the best configuration  $\mathbf{n}^*$  with the lowest  $q_c$  according to equation (2). The eigenvalue  $\lambda(\mathbf{n})$  (from now on we omit q in  $\lambda(\mathbf{n};q) \equiv \lambda(\mathbf{n})$ , which is always kept fixed) determines the growth rate of an arbitrary vector  $\mathbf{w}_0$  with 2M entries after  $\ell$  iterations of the matrix

$$\begin{split} \hat{\mathcal{M}}: & |w_{\ell}(\textbf{\textit{n}})| = \langle w_{\ell}|w_{\ell}\rangle^{\frac{1}{2}} = |\hat{\mathcal{M}}^{\ell}w_{0}| = \left\langle w_{0} \left| (\hat{\mathcal{M}}^{\ell})^{\dagger} \hat{\mathcal{M}}^{\ell} \right| w_{0} \right\rangle^{\frac{1}{2}} \sim \mathrm{e}^{\ell \log \lambda(\textbf{\textit{n}})}. \end{split}$$
 The largest eigenvalue is then calculated by the power method:

$$\lambda(\mathbf{n}) = \lim_{\ell \to \infty} \left[ \frac{|\mathbf{w}_{\ell}(\mathbf{n})|}{|\mathbf{w}_{0}|} \right]^{1/\ell} \tag{3}$$

Equation (3) is the starting point of an (infinite) perturbation series that provides the exact solution to the many-body influence problem in random networks and therefore contains all physical effects, including the collective influence. In practice, we minimize the cost energy function of influence  $|w_{\ell}(n)|$  in equation (3) for a finite  $\ell$ . The solution rapidly converges to the exact value as  $\ell \to \infty$ , the faster the larger the spectral gap. We find for  $\ell \ge 1$ , to leading order in 1/N (Supplementary Information section IIE):

$$|\mathbf{w}_{\ell}(\mathbf{n})|^2 = \sum_{i=1}^{N} (k_i - 1) \sum_{j \in \partial \text{Ball}(i, 2\ell - 1)} \left( \prod_{k \in \mathcal{P}_{2\ell - 1}(i, j)} n_k \right) (k_j - 1)$$
 (4)

where  $\mathrm{Ball}(i,\ell)$  is the set of nodes inside a ball of radius  $\ell$  (defined as the shortest path) around node  $i,\partial \mathrm{Ball}(i,\ell)$  is the frontier of the ball,  $\mathcal{P}_{\ell}(i,j)$  is the shortest path of length  $\ell$  connecting i and j (Fig. 1d), and  $k_i$  is the degree of node i.

The first collective optimization in equation (4) is  $\ell=1$ . We find  $|w_1(n)|^2=\sum_{i,j=1}^N A_{ij}(k_i-1)(k_j-1)n_in_j$ , where  $A_{ij}$  is the adjacency matrix (equation (39) in Supplementary Information). This term is interpreted as the energy of an antiferromagnetic Ising model with random bonds in a random external field at fixed magnetization, which is an example of a pair-wise NP-complete spin-glass whose solution is found in Supplementary Information section III with the cavity method<sup>28</sup> (Extended Data Fig. 2).

For  $\ell \geq 2$ , the problem can be mapped exactly to a statistical mechanical system with many-body interactions which can be recast in terms of a diagrammatic expansion, equations (41)–(49) in Supplementary Information. For example,  $|w_2(n)|^2$  leads to 4-body interactions (equation (45) in Supplementary Information), and, in general, the energy cost  $|w_\ell(n)|^2$  contains  $2\ell$ -body interactions. As soon as  $\ell \geq 2$ , the cavity method becomes much more complicated to implement and we use another suitable method, called extremal optimization (EO)<sup>29</sup> (Supplementary Information section IV). This method estimates the true optimal value of the threshold by finite-size scaling following extrapolation to  $\ell \to \infty$  (Extended Data Figs 3, 4). However, EO is not scalable to find the optimal configuration in large networks. Therefore, we develop an adaptive method, which performs excellently in practice, preserves the features of EO, and is highly scalable to present-day big data.

The idea is to remove the nodes causing the biggest drop in the energy function, equation (4). First, we define a ball of radius  $\ell$  around every node (Fig. 1d). Then, we consider the nodes belonging to the frontier  $\partial \text{Ball}(i,\ell)$  and assign to node i the collective influence (CI) strength at level  $\ell$  following equation (4):

$$\operatorname{CI}_{\ell}(i) = (k_i - 1) \sum_{j \in \partial \operatorname{Ball}(i, \ell)} (k_j - 1)$$
(5)

We notice that, while equation (4) is valid only for odd radii of the ball,  $\operatorname{CI}_\ell(i)$  is defined also for even radii. This generalization is possible by considering an energy function for even radii analogous to equation (4), as explained in Supplementary Information section IIG. The case of one-body interaction with zero radius  $\ell=0$  (equation (59) in Supplementary Information) leads to the high-degree (HD) ranking (equation (62) in Supplementary Information)<sup>10</sup>.

The collective influence, equation (5), is our second and most important result since it is the basis for the highly scalable and optimized CI algorithm which follows. In the beginning, all the nodes are present:  $n_i = 1$  for all i. Then, we remove node  $i^*$  with highest  $\text{CI}_\ell$  and set  $n_{i^*} = 0$ . The degree of each neighbour of  $i^*$  is decreased by one, and the procedure is repeated to find the new top CI node to remove. The algorithm is terminated when the giant component is zero (see Supplementary Information section V for implementation, and

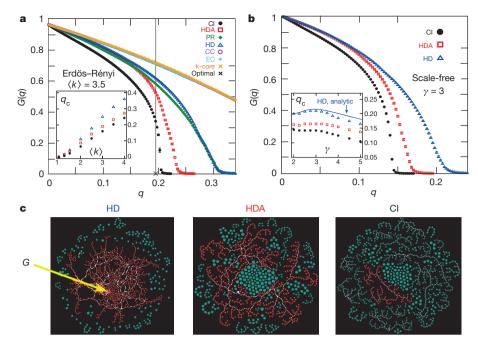


Figure 2 | Exact optimal solution and performance of CI in synthetic networks. a, G(q) in an ER network ( $N=2\times 10^5$ ,  $\langle k\rangle=3.5$ , error bars are s.e.m. over 20 realizations). We show the true optimal solution found with EO ('×' symbol), and also using CI, HDA, PR, HD, CC, EC and k-core methods. The other methods are not scalable and perform worse than HDA and are treated in Supplementary Information sections VI and VII (Extended Data Figs 8, 9). CI is close to the optimal  $q_c^{\rm opt}=0.192(9)$  obtained with EO in Supplementary Information section IV. Note that EO can estimate the extrapolated optimal value of  $q_c$ , but it cannot provide the optimal

Supplementary Information section VA for minimizing  $G(q) \neq 0$ ). By increasing the radius  $\ell$  of the ball we obtain better and better approximations of the optimal exact solution as  $\ell \to \infty$  (for finite networks,  $\ell$  does not exceed the network diameter).

The collective influence  $CI_{\ell}$  for  $\ell \geq 1$  has a rich topological content, and consequently tells us more about the role played by nodes in the network than the non-interacting high-degree hub-removal strategy at  $\ell = 0$ , CI<sub>0</sub>. The augmented information comes from the sum in the right hand side of equation (5), which is absent in the naive highdegree rank. This sum contains the contribution of the nodes living on the surface of the ball surrounding the central vertex i, each node weighted by the factor  $k_i - 1$ . This means that a node placed at the centre of a corona irradiating many links—the structure hierarchically emerging at different  $\ell$  levels as seen in Fig. 1e—can have a very large collective influence, even if it has a moderate or low degree. Such 'weak nodes' can outrank nodes with larger degree that occupy mediocre peripheral locations in the network. The commonly used word 'weak' in this context sounds particularly paradoxical. It is, indeed, usually used as a synonym for a low-degree node with an additional bridging property, which has resisted a quantitative formulation. We provide this definition through equation (5), according to which weak nodes are, de facto, quite strong. Paraphrasing Granovetter's conundrum<sup>30</sup>, equation (5) quantifies the "strength of weak nodes".

The CI-algorithm scales as  $\sim \mathcal{O}(N \log N)$  by removing a finite fraction of nodes at each step (Supplementary Information section VB). This high scalability allows us to find top influencers in current big-data social media and the minimal set of people to immunize in large-scale populations at the country level. The applications are investigated next.

Figure 2a shows the optimal threshold  $q_c$  for a random Erdös–Rényi (ER) network<sup>5</sup> (marked by the vertical line) obtained by extrapolating the EO solution to  $N \rightarrow \infty$  and  $\ell \rightarrow \infty$  (Supplementary Information section IV). In the same figure we compare the optimal threshold against the heuristic centrality measures: high-degree (HD)<sup>9</sup>, high-degree

configuration for large systems. Inset,  $q_c$  (obtained at the peak of the second-largest cluster) for the three best methods versus  $\langle k \rangle$ . **b**, G(q) for a SF network with degree exponent  $\gamma=3$ , maximum degree  $k_{\rm max}=10^3$ , minimum degree  $k_{\rm min}=2$  and  $N=2\times 10^5$  (error bars are s.e.m. over 20 realizations). Inset,  $q_c$  versus  $\gamma$ . The continuous blue line is the HD analytical result computed in Supplementary Information section IIG (Extended Data Fig. 1b). **c**, Example of SF network with  $\gamma=3$  after the removal of 15% of nodes, using the three methods HD, HDA and CI. CI produces a much reduced giant component G (red nodes)

adaptive (HDA), PageRank (PR)<sup>7</sup>, closeness centrality (CC)<sup>6</sup>, eigenvector centrality (EC)<sup>6</sup>, and k-core<sup>12</sup> (see Supplementary Information section I for definitions). Supplementary Information sections VI and VII show the comparison with the remaining heuristics<sup>6,11</sup> and the Belief Propagation method of ref. 14, respectively, which have worse computational complexity (and optimality), and cannot be applied to the network sizes used here. Remarkably, at the optimal value  $q_c$  predicted by our theory, the best among the heuristic methods (HDA, PR and HD) still predict a giant component  $\sim$ 50–60% of the whole original network. Furthermore, the influencer threshold predicted by CI approximates very well the optimal one, and, notably, CI outperforms the other strategies. Figure 2b compares CI in scale-free (SF) networks<sup>5</sup> against the best heuristic methods, that is, HDA and HD. In all cases, CI produces a smaller threshold and a smaller giant component (Fig. 2c).

As an example of an information spreading network, we consider the web of Twitter users (Supplementary Information section VIII<sup>19</sup>). Figure 3a shows the giant component of Twitter when a fraction q of its influencers is removed following CI. It is surprising that a lot of Twitter users with a large number of contacts have a mild influence on the network. This is witnessed by the fact that, when CI (at  $\ell = 5$ ) predicts a zero giant component (and so it exhausts the number of optimal influencers), the scalable heuristic ranks (HD, HDA, PR and k-core) still give a substantial giant component of the order of 30–70% of the entire network. These heuristics also, inevitably, find a remarkably large number of (fake) influencers, which is at least 50% larger than that predicted by CI (Fig. 3b and Supplementary Information section VIII). One cause for the poor performance of the high-degree-based ranks is that most of the hubs are clustered, which gives a mediocre importance to their contacts. As a consequence, hubs are outranked by nodes with lower degree surrounded by coronas of hubs (shown in detail in Fig. 3c), that is, the weak nodes predicted by the theory (Fig. 1e).

Finally, we simulate an immunization scheme on a personal contact network built from the phone calls performed by 14 million people in

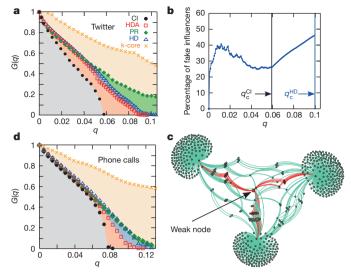


Figure 3 | Performance of CI in large-scale real social networks. a, Giant component G(q) of Twitter users<sup>19</sup> (N = 469,013) computed using CI, HDA, PR, HD and k-core strategies (other heuristics have prohibitive running times for this system size). b, Percentage of fake influencers or false positives (PFI, equation (120) in Supplementary Information) in Twitter as a function of q, defined as the percentage of non-optimal influencers identified by the HD algorithm in comparison with CI. Below  $q_c^{\text{CI}}$ , PFI reaches as much as ~40%, indicating the failure of HD in optimally finding the top influencers. Indeed, to obtain G = 0, HD has to remove a much larger number of fake influencers, which at  $q_c^{\rm HD}$  reaches PFI  $\approx 48\%$ . c, An example of the many weak nodes found in Twitter. These crucial influencers were missed by all heuristic strategies. **d**, G(q) for a social network of  $1.4 \times 10^7$  mobile phone users in Mexico representing an example of big data to test the scalability and performance of the algorithm in real networks. CI immunizes this social network using half a million fewer people than the best heuristic strategy (HDA), saving  $\sim$ 35% of the vaccine stockpile.

Mexico (Supplementary Information section IX). Figure 3d shows that our method saves a large number of vaccines or, equivalently, finds the smallest possible set of people to quarantine; our method therefore also outranks the scalable heuristics in large real networks. Thus, while the mapping of the influencer identification problem onto optimal percolation is strictly valid for locally tree-like random networks, our results may apply also to real loopy networks, provided the density of loops is not excessively large.

Our solution to the optimal influence problem shows its importance in that it helps to unveil hitherto hidden relations between people, as witnessed by the weak-node effect. This, in turn, is the by-product of a broader notion of influence, lifted from the individual non-interacting point of view<sup>6-12,19,20</sup> to the collective sphere: influence is an emergent property of collectivity, and top influencers arise from the optimization of the complex interactions they stipulate.

**Online Content** Methods, along with any additional Extended Data display items and Source Data, are available in the online version of the paper; references unique to these sections appear only in the online paper.

### Received 19 February; accepted 20 May 2015. Published online 1 July 2015.

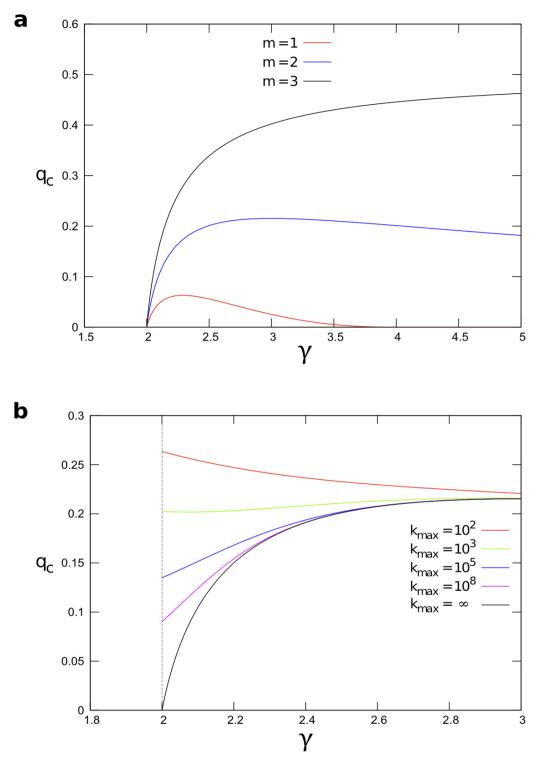
 Domingos, P. & Richardson, M. Mining knowledge-sharing sites for viral marketing. In Proc. 8th ACM SIGKDD Int. Conf. on Knowledge Discovery and Data Mining, 61–70 (ACM, 2002); http://dx.doi.org/10.1145/775047.775057.

- Pastor-Satorras, R. & Vespignani, A. Epidemic spreading in scale-free networks. Phys. Rev. Lett 86, 3200–3203 (2001).
- Newman, M. E. J. Spread of epidemic disease on networks. Phys. Rev. E 66, 016128 (2002)
- Kempe, D., Kleinberg, J. & Tardos, E. Maximizing the spread of influence through a social network. In *Proc. 9th ACM SIGKDD Int. Conf. on Knowledge Discovery* and *Data Mining*, 137–143 (ACM, 2003); http://dx.doi.org/10.1145/ 956750.956769.
- 5. Newman, M. E. J. Networks: An Introduction (Oxford Univ. Press, 2010).
- Freeman, L. C. Centrality in social networks: conceptual clarification. Soc. Networks 1, 215–239 (1978).
- Brin, S. & Page, L. The anatomy of a large-scale hypertextual web search engine. Comput. Networks ISDN Systems 30, 107–117 (1998).
- Kleinberg, J. Authoritative sources in a hyperlinked environment. In Proc. 9th ACM-SIAM Symp. on Discrete Algorithms (1998); J. Assoc. Comput. Machinery 46, 604–632 (1999).
- Albert, R., Jeong, H. & Barabási, A.-L. Error and attack tolerance of complex networks. Nature 406, 378–382 (2000).
- Cohen, R., Erez, K., ben-Avraham, D. & Havlin, S. Breakdown of the Internet under intentional attack. *Phys. Rev. Lett.* 86, 3682–3685 (2001).
- Chen, Y., Paul, G., Havlin, S., Liljeros, F. & Stanley, H. E. Finding a better immunization strategy. *Phys. Rev. Lett.* **101**, 058701 (2008).
- Kitsak, M. et al. Identification of influential spreaders in complex networks. Nature Phys. 6, 888–893 (2010).
- 13. Altarelli, F., Braunstein, A., Dall'Asta, L. & Zecchina, R. Optimizing spread dynamics on graphs by message passing. *J. Stat. Mech.* P09011 (2013).
- Altarelli, F., Braunstein, A., Dall'Asta, L., Wakeling, J. R. & Zecchina, R. Containing epidemic outbreaks by message-passing techniques. *Phys. Rev. X* 4, 021024 (2014)
- Hashimoto, K. Zeta functions of finite graphs and representations of p-adic groups. Adv. Stud. Pure Math. 15, 211–280 (1989).
- Coja-Oghlan, A., Mossel, E. & Vilenchik, D. A spectral approach to analyzing belief propagation for 3-coloring. Combin. Probab. Comput. 18, 881–912 (2009).
- Granovetter, M. Threshold models of collective behavior. Am. J. Sociol. 83, 1420–1443 (1978).
- Watts, D. J. A simple model of global cascades on random networks. Proc. Natl Acad. Sci. USA 99, 5766–5771 (2002).
- Pei, S., Muchnik, L., Andrade, J. S. Jr, Zheng, Z. & Makse, H. A. Searching for superspreaders of information in real-world social media. Sci. Rep. 4, 5547 (2014)
- Pei, S. & Makse, H. A. Spreading dynamics in complex networks. J. Stat. Mech. P12002 (2013).
- 21. Bollobás, B. & Riordan, O. *Percolation* (Cambridge Univ. Press, 2006).
- Bianconi, G. & Dorogovtsev, S. N. Multiple percolation transitions in a configuration model of network of networks. *Phys. Rev. E* 89, 062814 (2014).
- Karrer, B., Newman, M. E. J. & Zdeborová, L. Percolation on sparse networks. Phys. Rev. Lett. 113, 208702 (2014).
- 24. Angel, O., Friedman, J. & Hoory, S. The non-backtracking spectrum of the universal cover of a graph. *Trans. Am. Math. Soc.* **367**, 4287–4318 (2015).
- Krzakala, F. et al. Spectral redemption in clustering sparse networks. Proc. Natl Acad. Sci. USA 110, 20935–20940 (2013).
- Newman, M. E. J. Spectral methods for community detection and graph partitioning. *Phys. Rev. E* 88, 042822 (2013).
- Radicchi, F. Predicting percolation thresholds in networks. Phys. Rev. E 91, 010801(R) (2015).
- Mézard, M. & Parisi, G. The cavity method at zero temperature. J. Stat. Phys. 111, 1–34 (2003).
- 29. Boettcher, S. & Percus, A. G. Optimization with extremal dynamics. *Phys. Rev. Lett.* **86**, 5211–5214 (2001).
- 30. Granovetter, M. The strength of weak ties. Am. J. Sociol. 78, 1360–1380 (1973).

**Supplementary Information** is available in the online version of the paper.

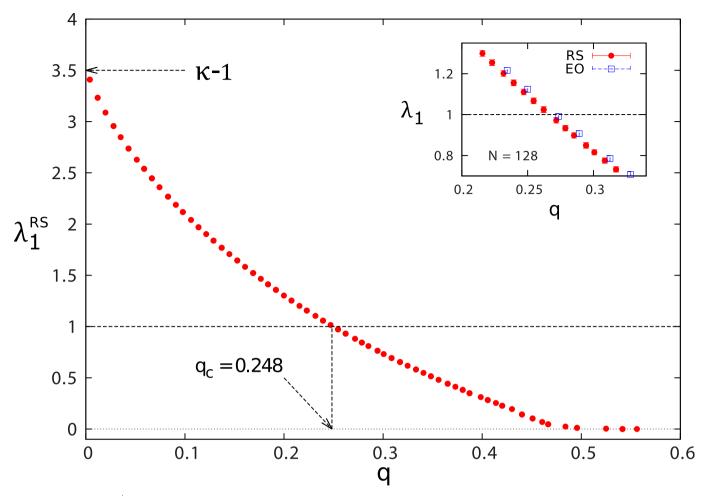
**Author Contributions** Both authors contributed equally to the work presented in this paper.

**Author Information** Reprints and permissions information is available at www.nature.com/reprints. The authors declare no competing financial interests. Readers are welcome to comment on the online version of the paper. Correspondence and requests for materials should be addressed to H.A.M. (hmakse@lev.ccny.cuny.edu).



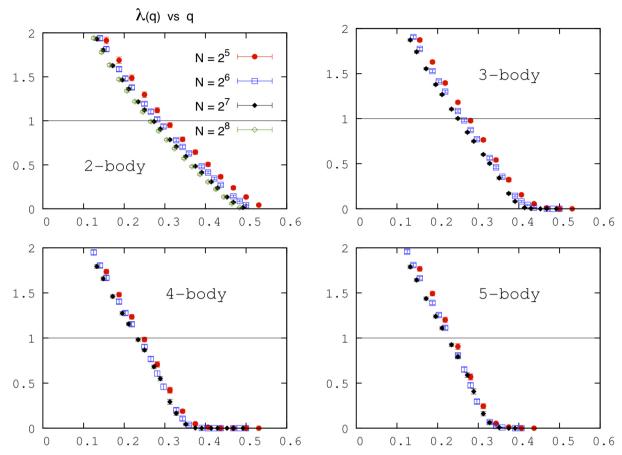
Extended Data Figure 1 | High-degree (HD) threshold. a, HD influence threshold  $q_c$  as a function of the degree distribution exponent  $\gamma$  of scale-free networks in the ensemble with  $k_{\max} = mN^{1/(\gamma-1)}$  and  $N \to \infty$ . The curves refer to different values of the minimum degree m: 1 (red), 2 (blue), 3 (black). The fragility of SF networks (small  $q_c$ ) is notable for m=1 (the case calculated in ref. 10). In this case (m=1), the network contains many leaves, and reduces to a star at  $\gamma=2$ , which is trivially destroyed by removing the only single hub, explaining the general fragility in this case. Furthermore, in this same case, the network becomes a collection of dimers with k=1 when  $\gamma\to\infty$ , which is still trivially fragile. This also explains why  $q_c\to 0$  for  $\gamma\ge 4$ . Therefore, the fragility in the case m=1 has its roots in these two limiting trivial cases. Removing the leaves (m=2) results in a 2-core, which is already more robust.

For the 3-core m=3,  $q_c\approx 0.4$ –0.5 provides a quite robust network, and has the expected asymptotic limit to a non-zero  $q_c$  of a random regular graph with k=3 as  $\gamma\to\infty$ ,  $q_c\to(k-2)/(k-1)=0.5$ . Thus, SF networks become robust in these more realistic cases, and the search for other attack strategies becomes even more important. **b**, HD influence threshold  $q_c$  as a function of the degree distribution exponent of scale-free networks with minimum degree m=2 in the ensemble where  $k_{\rm max}$  is fixed and does not scale with N. The curves refer to different values of the cut-off  $k_{\rm max}$ :  $10^2$  (red),  $10^3$  (green),  $10^5$  (blue),  $10^8$  (magenta), and  $k_{\rm max}=\infty$  (black), and show that for a typical  $k_{\rm max}$  degree of  $10^3$ , for instance in social networks, the network is fairly robust with  $q_c\approx 0.2$  for all  $\gamma$ . The curve with m=2 and  $k_{\rm max}=10^3$  is replotted in the inset of Fig. 2b.



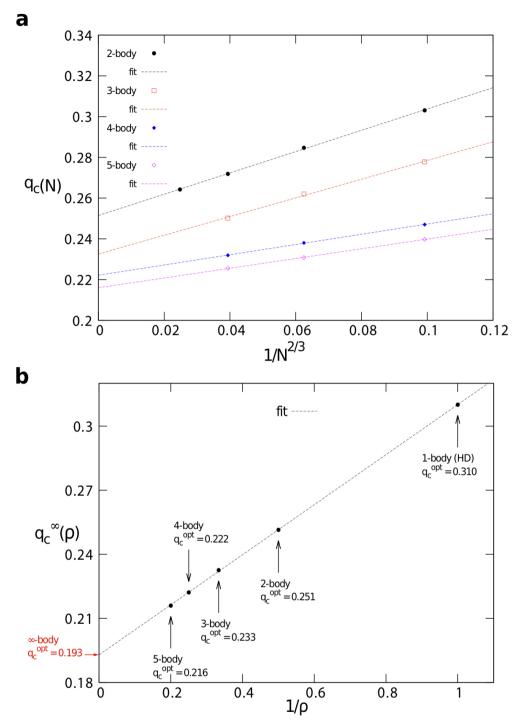
Extended Data Figure 2 | Replica Symmetry (RS) estimation of the maximum eigenvalue. Main panel, the eigenvalue  $\lambda_1^{RS}(q)$ , equation (92) in Supplementary Information for the two-body interaction  $\ell=1$ , obtained by minimizing the energy function  $\mathcal{E}(s)$  with the RS cavity method. The curve was computed on an ER graph of N=10,000 nodes and average degree  $\langle k \rangle=3.5$ 

and then averaged over 40 realizations of the network (error bars are s.e.m.). Inset, comparison between the RS cavity method and EO (extremal optimization) for an ER graph of  $\langle k \rangle = 3.5$  and N=128. The curves are averaged over 200 realizations (error bars are s.e.m.).



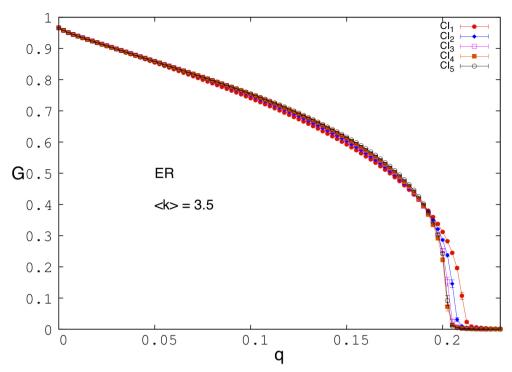
Extended Data Figure 3 | EO estimation of the maximum eigenvalue. Eigenvalue  $\lambda(q)$  obtained by minimizing the energy function  $\mathcal{E}(\mathbf{n})$  with  $\tau EO$  ( $\tau$ -extremal optimization), plotted as a function of the fraction of removed nodes q. The panels are for different orders of the interactions. The curves in

each panel refer to different sizes of ER networks with average connectivity  $\langle k \rangle = 3.5$ . Each curve is an average over 200 instances (error bars are s.e.m.). The value  $q_{\rm c}$  where  $\lambda(q_{\rm c})=1$  is the threshold for a particular N and many-body interaction.



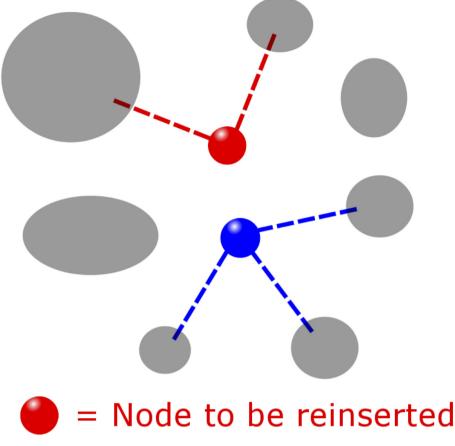
Extended Data Figure 4 | Estimation of optimal threshold  $q_c^{\rm opt}$  with EO. a, Critical threshold  $q_c$  as a function of the system size N, obtained with EO from Extended Data Fig. 3, of ER networks with  $\langle k \rangle = 3.5$  and varying size. The curves refer to different orders of the many-body interactions. The data show a linear behaviour as a function of  $N^{-2/3}$ , typical of spin glasses, for each many-

body interaction  $\rho$ . The extrapolated value  $q_c^\infty(\rho)$  is obtained at the y intercept. **b**, Thermodynamic critical threshold  $q_c^\infty(\rho)$  as a function of the order of the interactions  $\rho$  from **a**. The data scale linearly with  $1/\rho$ . From the y intercept of the linear fit we obtain the thermodynamic limit of the infinite-body optimal value  $q_c^{\text{opt}} = q_c^\infty(\rho \to \infty) = 0.192(9)$ .



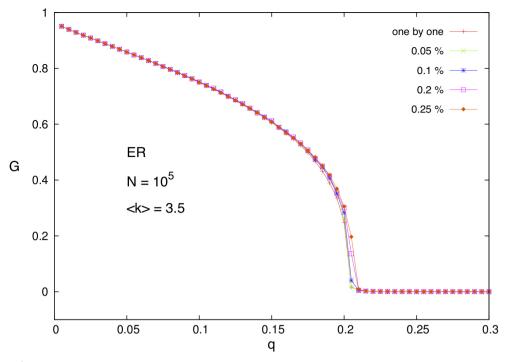
Extended Data Figure 5 | Comparison of the CI algorithm for different radii  $\ell$  of the Ball( $\ell$ ). We use  $\ell=1,2,3,4,5,$  on a ER graph with average degree  $\langle k \rangle = 3.5$  and  $N=10^5$  (the average is taken over 20 realizations of the network, error bars are s.e.m.). For  $\ell=3$  the performance is already practically indistinguishable from  $\ell=4$ , 5. The stability analysis we developed to minimize  $q_c$  is strictly valid only when G=0, since the largest eigenvalue of the

modified NB matrix controls the stability of the solution G=0, and not the stability of the solution G>0. In the region where G>0 we use a simple and fast procedure to minimize G explained in Supplementary Information section VA. This explains why there is a small dependence on having a slightly larger G for larger  $\ell$ , when G>0 in the region  $q\approx 0.15$ .

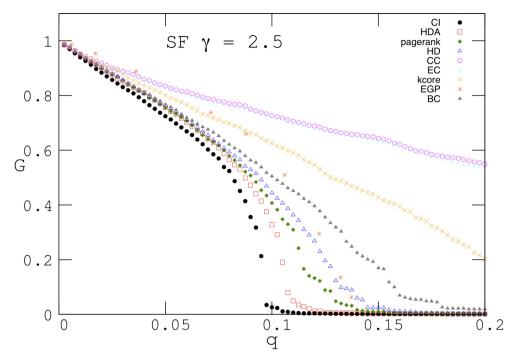


Extended Data Figure 6 | Illustration of the algorithm used to minimize G(q) for  $q < q_c$ . Starting from the completely fragmented network at  $q = q_c$ , the  $Nq_c$  influencers are reinserted with their original degree and connected to their original neighbours with the following criterion: each node is assigned and index c(i) given by the number of clusters it would join if it were reinserted

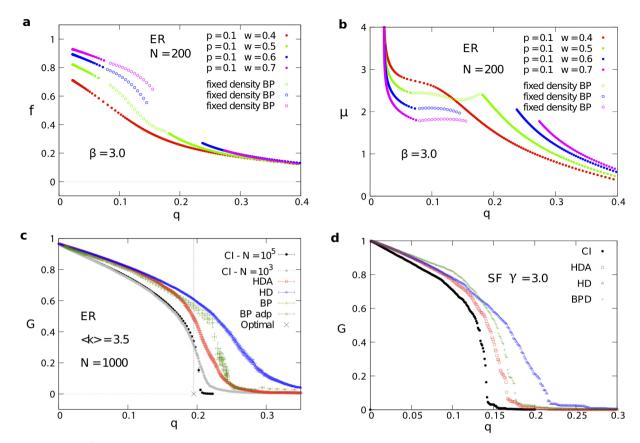
in the network. For example, the red node has c(red) = 2, while the blue one has c(blue) = 3. The node with the smallest c(i) is reinserted in the network: in this case the red node. Then the c(i)s are recalculated and the new node with the smallest c(i) is found and reinserted. These steps are repeated until all the removed nodes are reinserted in the network.



**Extended Data Figure** 7 | **Test of the decimation fraction.** Giant component *G* as a function of the fraction of removed nodes *q* using CI, for an ER network of  $N = 10^5$  nodes and average degree  $\langle k \rangle = 3.5$ . The profiles of the curves are drawn for different percentages of nodes fixed at each step of the decimation algorithm.

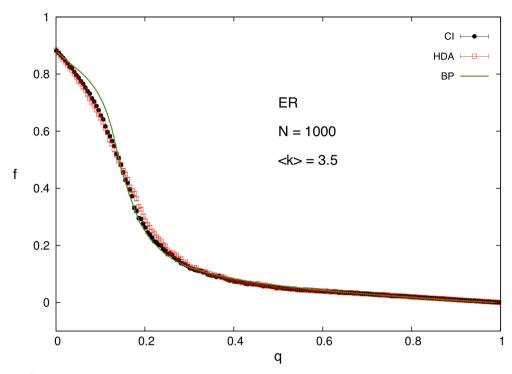


**Extended Data Figure 8** | **Comparison of the performance of CI, BC and EGP in destroying** *G***.** We also include HD, HDA, EC, CC, k-core and PR. We use a scale-free (SF) network with degree exponent  $\gamma = 2.5$ , average degree  $\langle k \rangle = 4.68$ , and  $N = 10^4$ . We use the same parameters as in ref. 11.



**Extended Data Figure 9** | **Comparison with BP for a network immunization. a**, Fraction of infected nodes f as a function of the fraction of immunized nodes q in the susceptible-infected-removed (SIR) model from the BP solution. We use an ER random graph of N=200 nodes and average degree  $\langle k \rangle = 3.5$ . The fraction of initially infected nodes is p=0.1 and the inverse temperature  $\beta=3.0$ . The profiles are drawn for different values of the transmission probability w: 0.4 (red curve), 0.5 (green), 0.6 (blue), 0.7 (magenta). Also shown are the results of the fixed density BP algorithm (open circles). **b**, Chemical potential  $\mu$  as a function of the immunized nodes q from BP. We use an ER random graph of N=200 nodes and average degree  $\langle k \rangle = 3.5$ . The fraction of the initially infected nodes is p=0.1 and the

inverse temperature  $\beta=3.0$ . The profiles are drawn for different values of the transmission probability w: 0.4 (red curve), 0.5 (green), 0.6 (blue), 0.7 (magenta). Also shown are the results of the fixed density BP algorithm (open circles) for the region where the chemical potential is non-convex. **c**, Comparison between the giant components obtained with CI, HDA, HD and BP. We use an ER network of  $N=10^3$  and  $\langle k \rangle=3.5$ . We also show the solution of CI from Fig. 2a for  $N=10^5$ . We find in order of performance: CI, HDA, BP and HD. (The average is taken over 20 realizations of the network, error bars are s.e.m.) **d**, Comparison between the giant components obtained with CI, HDA, HD and BPD. We use a SF network with degree exponent  $\gamma=3.0$ , minimum degree  $k_{\min}=2$ , and  $N=10^4$  nodes.



Extended Data Figure 10 | Fraction of infected nodes f(q) as a function of the fraction of immunized nodes q in SIR from BP. We use the following parameters: initial fraction of infected people p=0.1, and transmission probability w=0.5. We use an ER network of  $N=10^3$  nodes and  $\langle k \rangle = 3.5$ .

We compare CI, HDA and BP. All strategies give similar performance, owing to the large value of the initial infection p, which washes out the optimization performed by any sensible strategy, in agreement with the results shown in figure 12a of ref. 14.